

Controllability analysis of the magnetic flux distribution in ITER Hybrid scenarios

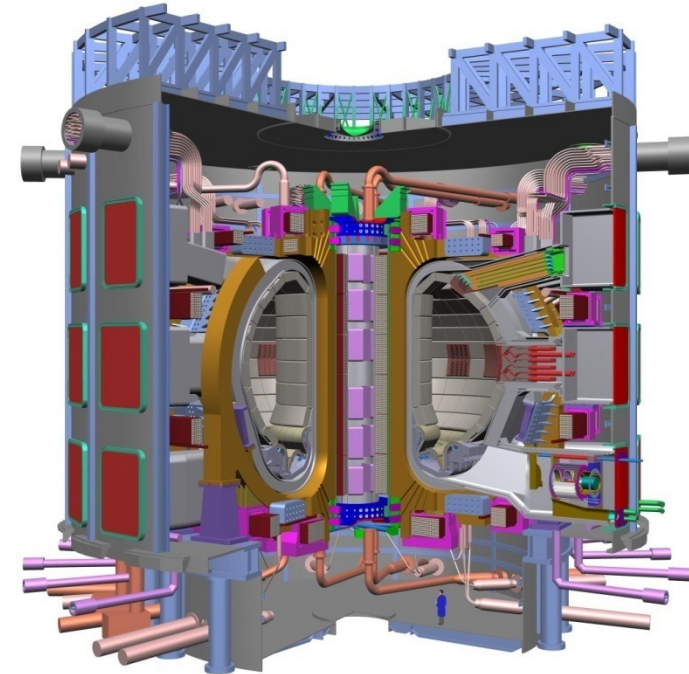
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ITER Goals for hybrid' scenario (long discharge time)

- $P_{\text{fus}} > 350\text{MW}$
- $I_p = 11\text{-}13\text{ MA}$
- $P_{\text{fus}} > 5 \cdot P_{\text{in}}$
- $t_{\text{discharge}} > 1000\text{ s}$
- $q > 1$ (for stability)



ITER hybrid scenario primary mission is to maximize neutron fluence per pulse, for reactor-relevant component testing: **Long reliable discharge with sufficient neutron flux**

CRONOS simulations carried out with the 'GLF23' turbulent transport model

s/q effect on transport tested by simulating various current drive mixes

(J.Citrin, G.M.D.Hogeweij, J.Garcia, J.F. Artaud, F.Imbeaux, Nucl. Fusion 2010)



Importance of q-profile shaping



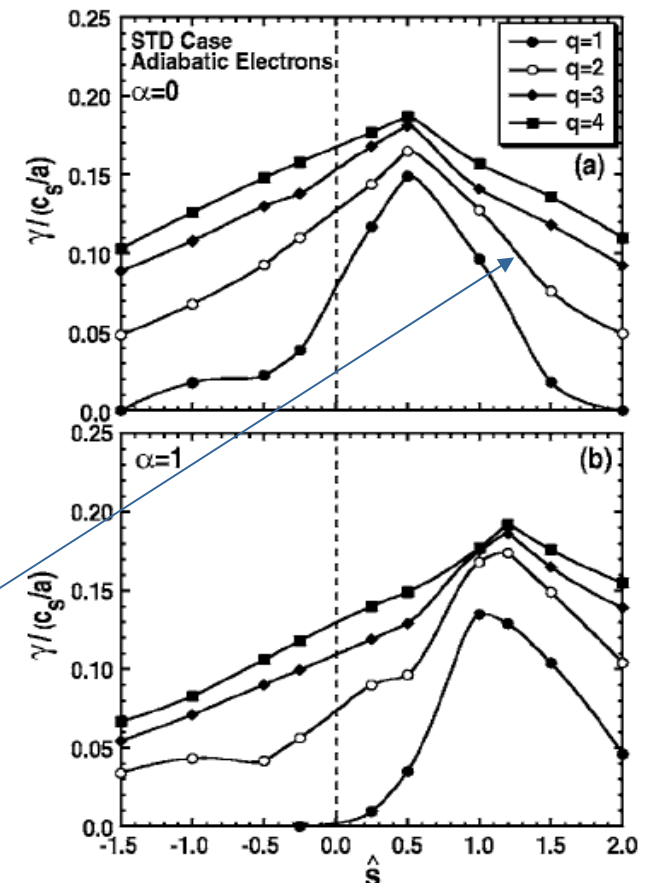
q-profile shape can be **optimized for improved confinement**. Theoretical models– including the GLF23 transport model – predict instability thresholds linearly dependent on s/q (for $s > \sim 0.5$)

GYRO ITG linear growth rates

Optimizing the q-profile shape with non-inductive current drive reduces the density necessary for $P_{fus}=350\text{MW}$.

↓

At the same Greenwald fraction, the current can thus also be reduced, **decreasing I_{ohm}** , and **increasing $t(q=1)$** .



In the moderate/high magnetic shear regime, we want to maximize s/q .

Kinsey, Waltz, Candy (Phys. Plasmas 2006)



Control oriented modeling



The poloidal magnetic flux $\psi(x, t)$ at any point in the poloidal cross section is the total flux through the surface S bounded by the toroidal ring passing through the point, i.e.,

$$\psi(x, t) = \frac{1}{2\pi} \int B_{\text{pol}} dS.$$

A control-orientated model of the magnetic flux

$$\frac{\partial \psi(x, t)}{\partial t} = \frac{\eta_{\parallel}(x)}{\mu_0 a_e^2} \left(\frac{\partial^2 \psi(x, t)}{\partial x^2} + \frac{1}{x} \frac{\partial \psi(x, t)}{\partial x} \right) + \eta_{\parallel}(x) R_0 j_{\text{ni}}(x, t),$$

where $j_{\text{ni}}(x, t) = j_{\text{bc}}(x) + j_{\text{nbi}}(x) + j_{\text{eccd}}(x, t)$ with the boundary conditions

$$\frac{\partial \psi(0, t)}{\partial x} = 0 \quad \frac{\partial \psi(1, t)}{\partial x} = -\frac{R_0 \mu_0 I_p}{2\pi}$$



Discretization distributed model



Due to the fact that the parameters $\eta_{\parallel}(x)$, $j_{bc}(x)$, and $j_{nbi}(x)$ are space dependent parameters, the PDE model has to be discretized in order to evaluate the magnetic flux $\psi(x, t)$.

$$\begin{aligned} \frac{d\psi(x_i, t)}{dt} = & \frac{\eta_{\parallel ij}}{\mu_0 a_e^2} (c_1(i)\psi_{i+1} - c_2(i)\psi_i + c_3(i)\psi_{i-1}) \\ & + \eta_{\parallel ij} R_0 (j_{bc}(x_i) + j_{nbi}(x_i) + j_{eccd}(x_i, t)) \end{aligned}$$

with the discretization coefficients

$$c_1(i) = 1/2 \frac{2x_i + \delta x}{\delta x^2 x_i}$$

$$c_2(i) = 2 \frac{1}{\delta x^2}$$

$$c_3(i) = 1/2 \frac{2x_i + \delta x}{\delta x^2 x_i}$$



The state-space form of the spatially discretized model

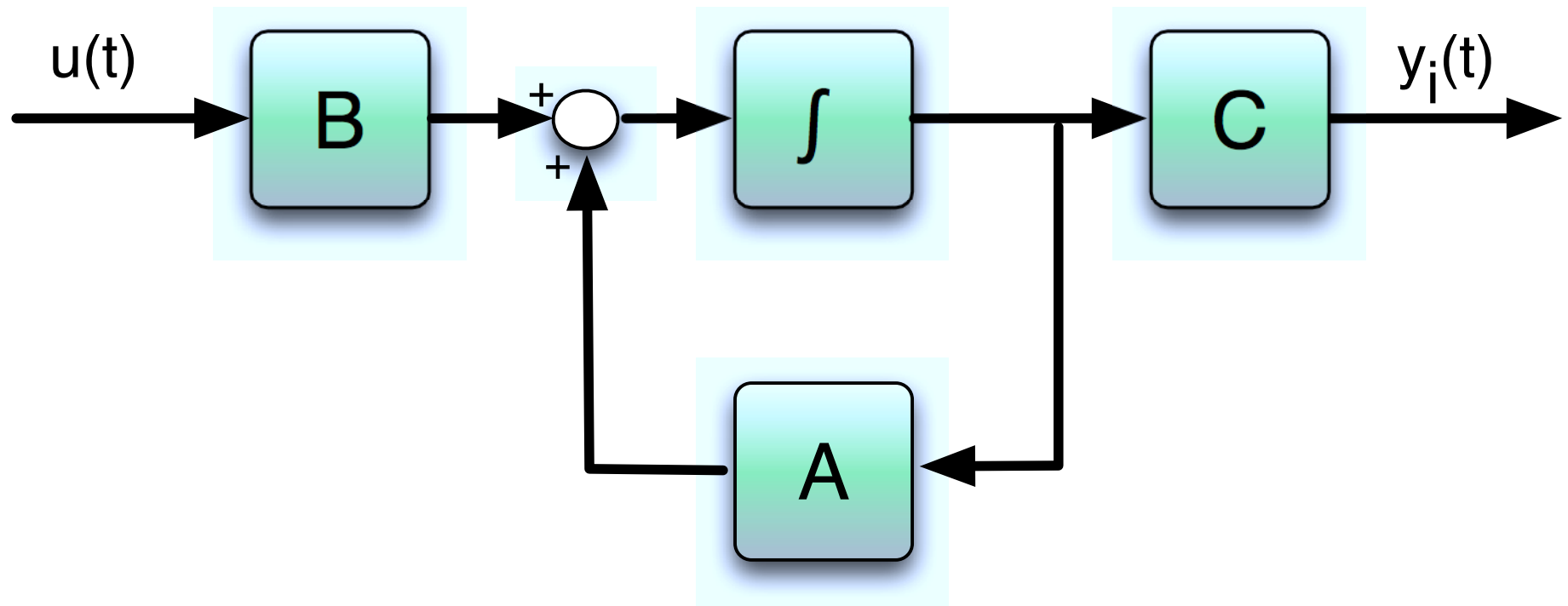
$$\begin{aligned}\frac{d\psi_i(t)}{dt} &= \mathbf{A}\psi_i(t) + \mathbf{B}u_i(t) \\ y_i(t) &= \mathbf{C}\psi_i(t)\end{aligned}$$

$\psi_i(t)$:	state vector	\mathbb{R}^N
$y_i(t)$:	output vector	\mathbb{R}^M
$u_i(t)$:	input vector	\mathbb{R}^R
\mathbf{A} :	system matrix	$\mathbb{R}^{N \times N}$
\mathbf{B} :	input matrix	$\mathbb{R}^{N \times R}$
\mathbf{C} :	output matrix	$\mathbb{R}^{M \times N}$

Transfer function: $\mathbf{G}(s) = \frac{y_i(s)}{u_i(s)} = \mathbf{C}(s\mathbf{I} - \mathbf{A})^{-1}\mathbf{B}$

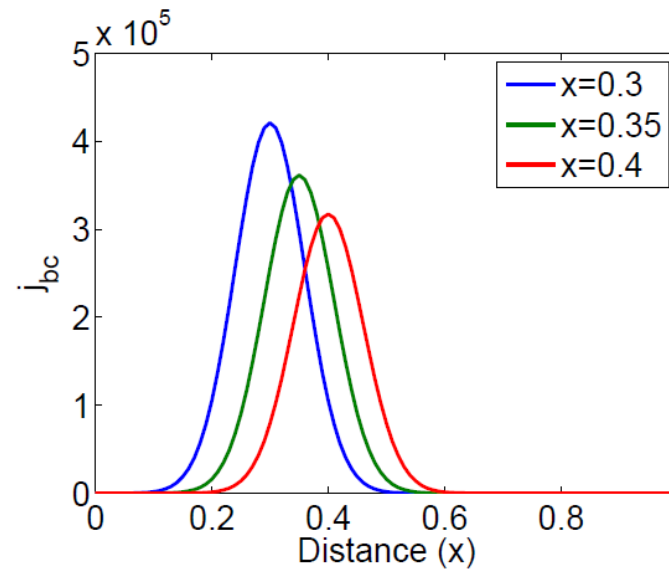
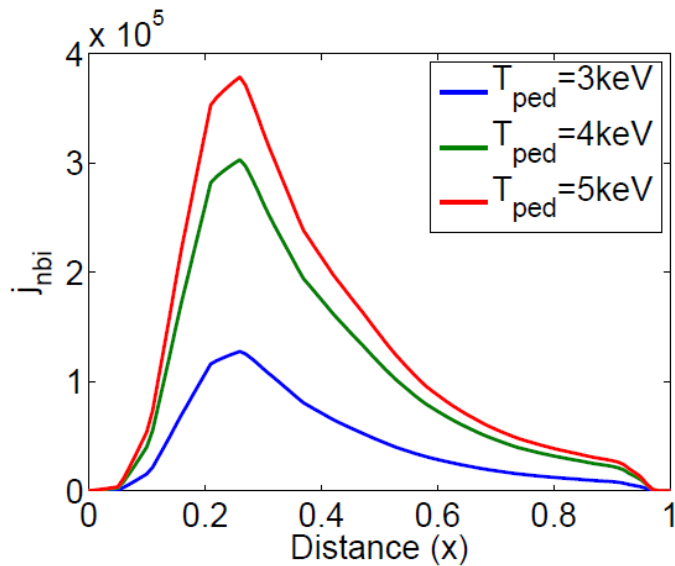
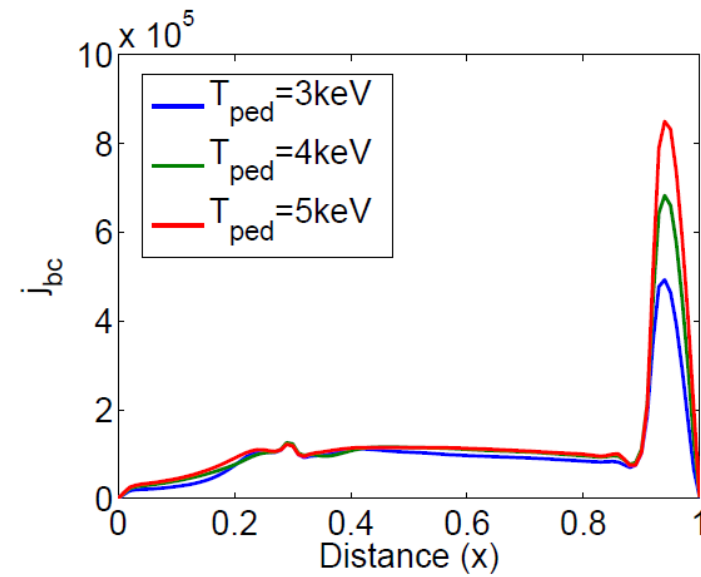
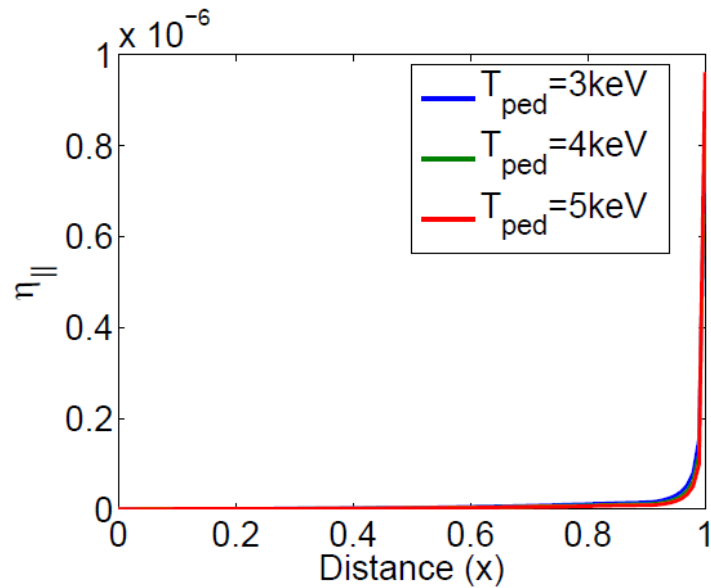


Graphical representation of State-Space model



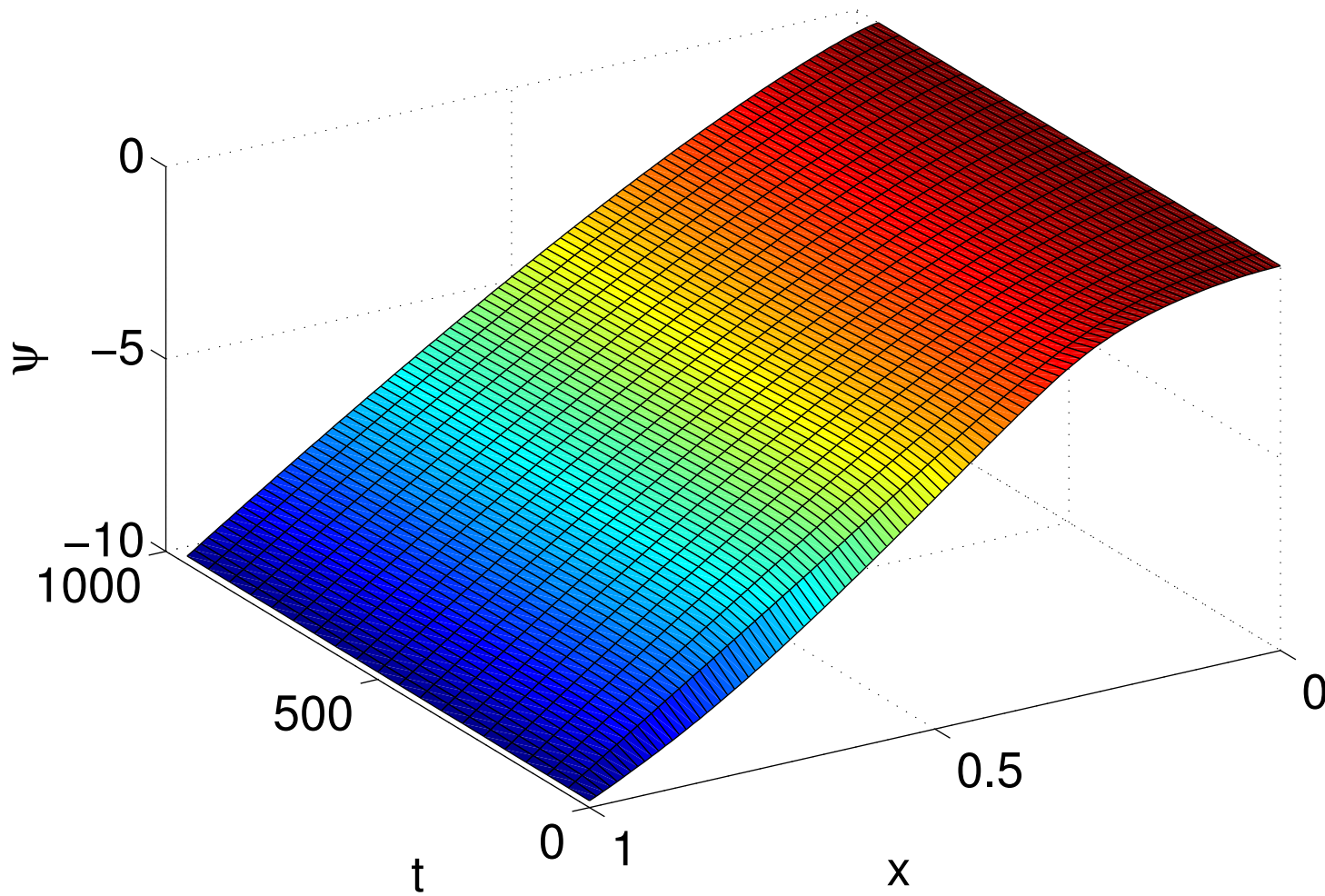


Input to state-space model from CRONOS





Variation poloidal flux with state-space model





Controllability analysis



The controllability matrix \mathcal{C}_i equals

$$\mathcal{C}_i(\mathbf{A}, \mathbf{B}) = [\mathbf{B} \quad \mathbf{AB} \quad \mathbf{A}^2\mathbf{B} \quad \dots \quad \mathbf{A}^{i-1}\mathbf{B}]$$

The observability matrix $\mathcal{O}_i(\mathbf{A}, \mathbf{C})$

$$\mathcal{O}_i(\mathbf{A}, \mathbf{C}) = \begin{bmatrix} \mathbf{C} \\ \mathbf{CA} \\ \mathbf{CA}^2 \\ \vdots \\ \mathbf{CA}^{i-1} \end{bmatrix}$$

The system is controllable/observable if the controllability/observability matrix has full rank.

$$\text{rank}(\mathcal{C}_i) = N \quad \text{rank}(\mathcal{O}_i) = N$$



Controllability / observability Gramian and reachable sets



For large scale systems, where $N > 10$, there are elements that require a significant amount of energy in terms of states.

The controllability/observability Gramian

$$\mathcal{P} = \mathcal{C}_\infty(\mathbf{A}, \mathbf{B})\mathcal{C}_\infty^T(\mathbf{A}, \mathbf{B}) = \sum_{i=0}^{\infty} \mathbf{A}^i \mathbf{B} \mathbf{B}^T (\mathbf{A}^T)^i,$$

$$\mathcal{Q} = \mathcal{O}_\infty(\mathbf{C}, \mathbf{A})\mathcal{O}_\infty^T(\mathbf{C}, \mathbf{A}) = \sum_{i=0}^{\infty} (\mathbf{A}^T)^i \mathbf{C}^T \mathbf{C} \mathbf{A}^i$$

The reachable sets from the given initial condition

$$J_{\text{con}}(\psi_i) = \psi_i^T \mathcal{P}^{-1} \psi_i \quad J_{\text{obs}}(\psi_i) = \psi_i^T \mathcal{Q} \psi_i$$



Singular values for accessible states



In general, a state coordinate transformation produces an equivalent model in another coordinate system in which $\bar{\psi}_i(t) = \mathbf{T}\psi_i(t)$

$$\begin{aligned}\frac{d\bar{\psi}_i(t)}{dt} &= \mathbf{TAT}^{-1}\bar{\psi}_i(t) + \mathbf{TB}u_i, \\ y_i &= \mathbf{CT}^{-1}\bar{\psi}_i(t),\end{aligned}$$

The associated Gramians $\bar{\mathcal{P}}$ and $\bar{\mathcal{Q}}$

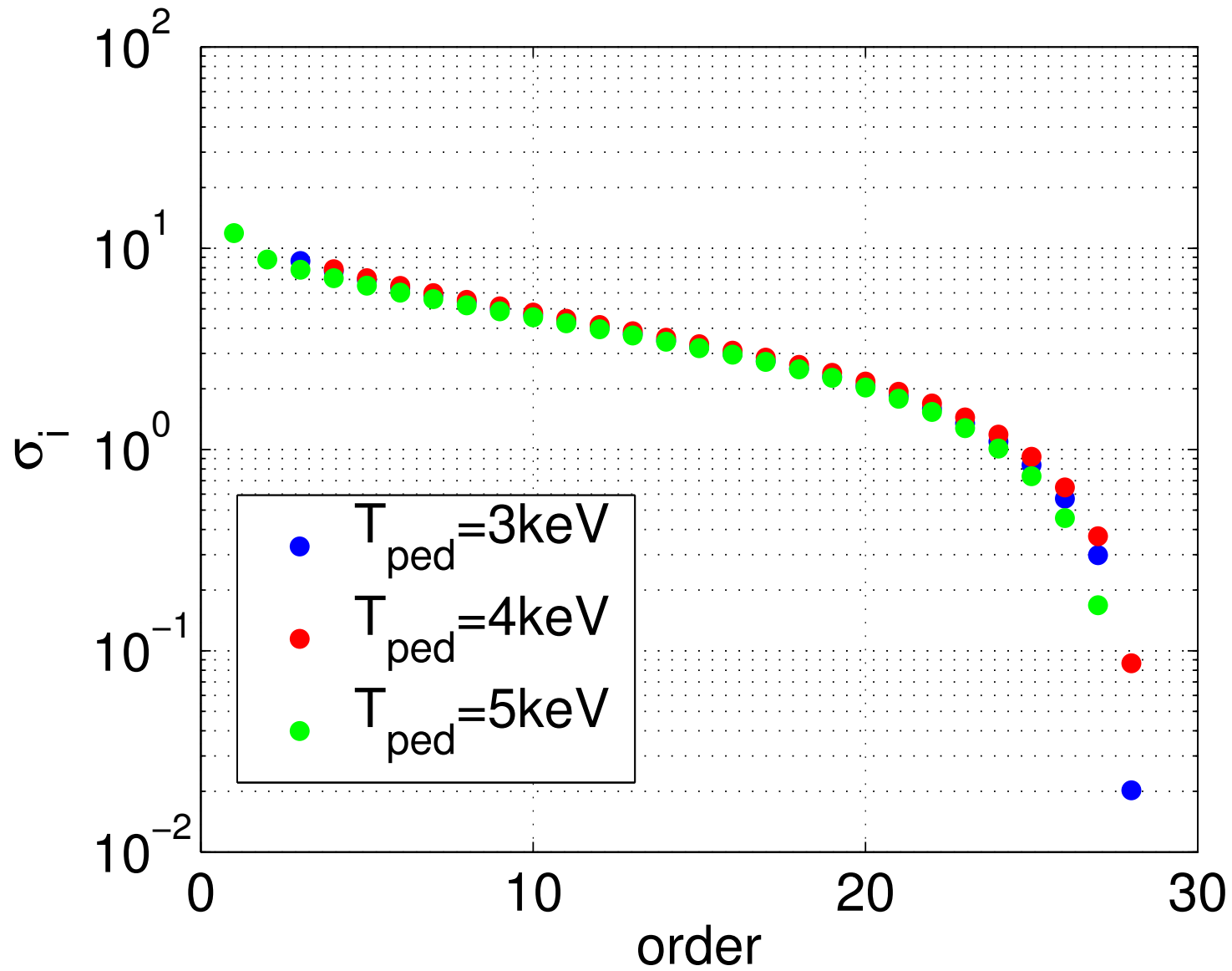
$$\bar{\mathcal{P}} = \mathbf{TP}\mathbf{T}^T \quad \bar{\mathcal{Q}} = \mathbf{T}^{-T}\mathcal{Q}\mathbf{T}^{-1} \quad \bar{\mathcal{P}}\bar{\mathcal{Q}} = \mathbf{TP}\mathcal{Q}\mathbf{T}^{-1}$$

$$\bar{\mathcal{P}} = \bar{\mathcal{Q}} = \text{diag}(\sigma_1, \sigma_2, \dots, \sigma_N)$$

$$\sigma_i = \sqrt{\lambda_i(\mathcal{P}\mathcal{Q})} = \sqrt{\lambda_i(\bar{\mathcal{P}}\bar{\mathcal{Q}})}, \quad i = 1, 2, \dots, N$$

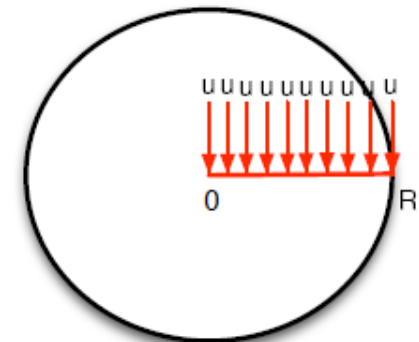
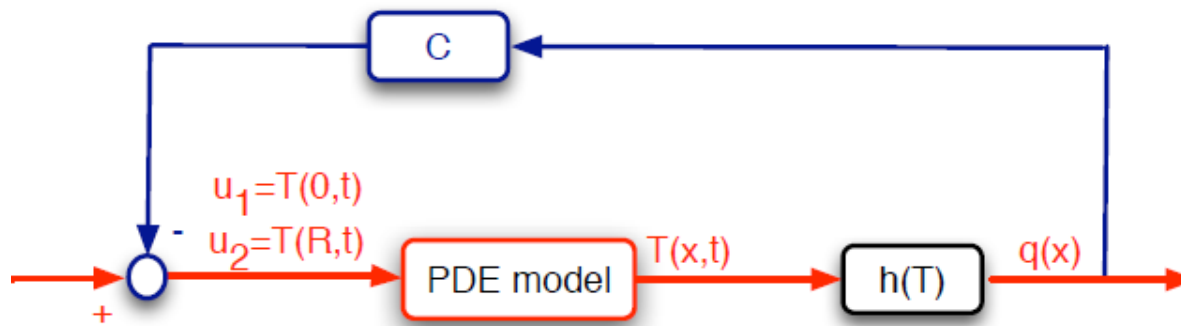
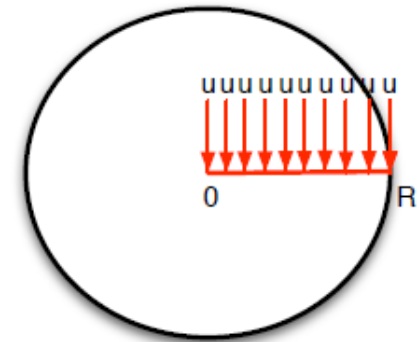
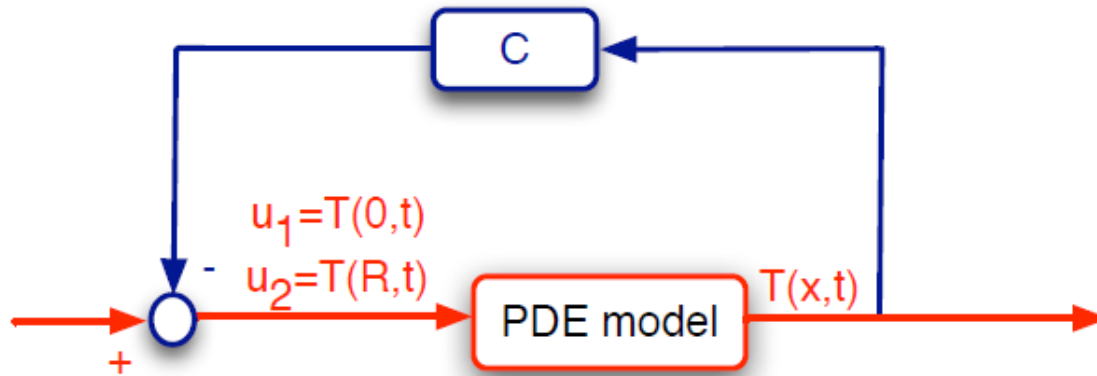


Singular values





Control





Summary



- Energy confinement in tokamaks limited by turbulence driven by temperature gradients
- Turbulence onset characterized by critical gradients, predicted by linear theory
- Critical gradients predicted to be sensitive to the current profile
- Discretized distributed state-space model for the flux distribution in ITER-Hybrids derived
- Observability / Controlability analysis carried out for ITER Hybrid scenarios, using detailed input from Cronos simulations
- Simulation suggests that only a subset of the states is effectively accessible
- Model based closed-loop control optimization remains to be carried out (before October 2011) based on $\langle s/q \rangle$ cost function. Expected to yield overall improved q and s -profiles (Also maximization of low s volume)