

Integrated core- SOL-divertor simulations of ITER H-mode scenarios with different pedestal density

I. Ivanova-Stanik^a, F. Koechl^b, I. Voitsekhovitch^c, A. Polevoi^d,
G. Telesca^e, R. Zagorski^a,
and the EU-ITM ITER Scenario Modelling group

^a*Institute of Plasma Physics and Laser Microfusion, EURATOM/IPPLM Association, Warsaw, Poland*

^b*Association EURATOM/OAW, Atominstitut, TU Wien, Vienna, Austria*

^c*EURATOM/CCFE Association, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK*

^d*ITER Organization, Route de Vinon sur Verdon, 13115 St Paul Lez Durance, France*

^e*Department of Applied Physics, Ghent University, B-9000 Gent, Belgium*

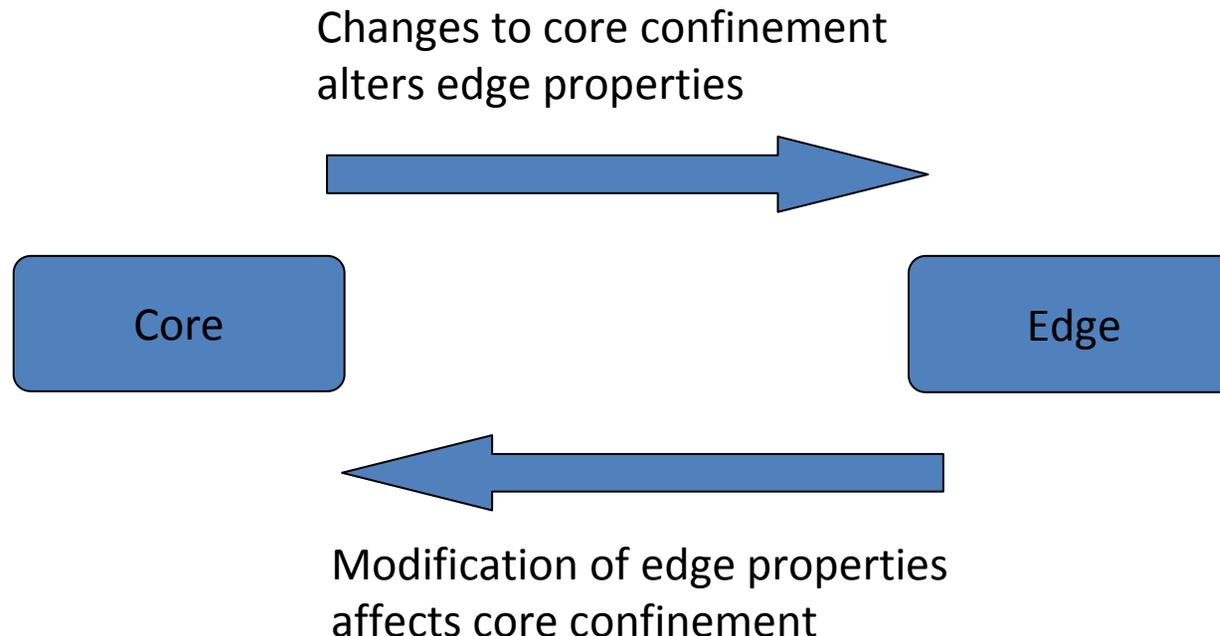
4th Int. Workshop on Plasma Edge Theory in Fusion Devices, 23-25 September 2013

COREDIV = 1D transport in the core self-consistently coupled to 2D model in the SOL

INTERPLAY BETWEEN EDGE AND CORE

Interaction between bulk and edge plasma is a complex phenomena involving a lot of physics processes mostly through:

- **Energy balance**
- **Particle balance**



Equations and transport model for the core plasma - Background ions

1D transport of particles (n_i) and energy (T_e , T_i):

$$\frac{\partial n_i}{\partial t} + \frac{1}{rg_1} \frac{\partial}{\partial r} \left[rg_2 \left(-D_i \frac{\partial n_i}{\partial r} + w_i n_i \right) \right] = S_i(r) = S_i^0 \times P(r) + \text{neocl. contrib.}$$

$$\frac{3}{2} \frac{\partial n_i T_i}{\partial t} + \frac{1}{rg_1} \frac{\partial}{\partial r} \left[rg_2 \left(-k_i \frac{\partial T_i}{\partial r} + \frac{5}{2} \Gamma_i T_i \right) \right] = P_{AUX}^i + Q_{ei}$$

S_i^0 iterated to have
constant $\langle n_e \rangle$

$$\frac{3}{2} \frac{\partial n_e T_e}{\partial t} + \frac{1}{rg_1} \frac{\partial}{\partial r} \left[rg_2 \left(-k_e \frac{\partial T_e}{\partial r} + \frac{5}{2} \Gamma_e T_e \right) \right] = P_{OH} + P_{AUX}^e + P_\alpha - P_B - P_{cyc} - P_{lin} - P_{ion} - Q_{ei}$$

Quasineutrality: $n_e = n_i + \sum_{k,j} n_j^k \quad \langle n_e \rangle = \text{const.} \quad i=D,T$

Equations for plasma rotation
and plasma current neglected \rightarrow
 $j(r)$ – given input function

g_1, g_2 – metric coefficients

Transport model for background ions

Anomalous transport described by simple model:

τ_E from experimental ELMy H-mode scaling

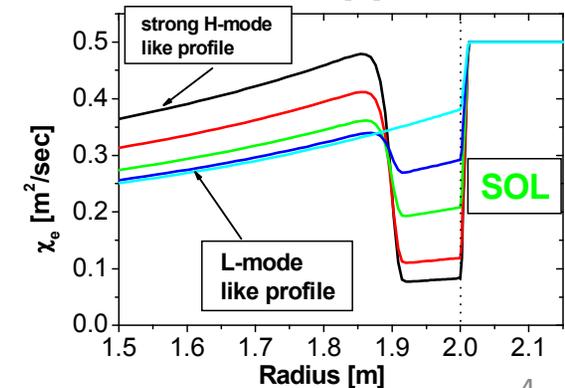
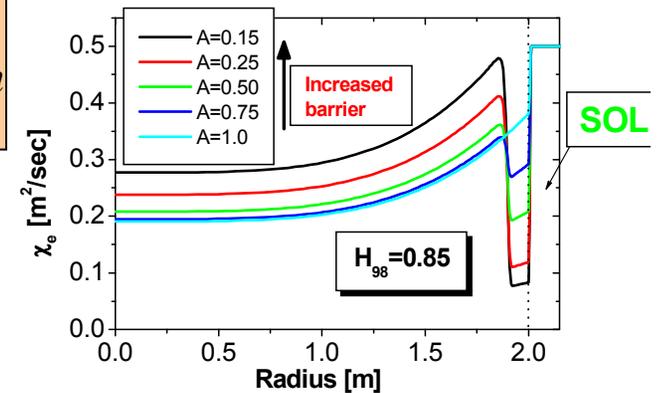
$$\chi_e^{an} = C_e \frac{a^2}{\tau_E} F(r)$$

$$\chi_i^{an} = \chi_e^{an} \quad D_i^{an} = 0.35 \chi_e^{an}$$

F(r) – profile function

$$F(r) = \left(1 + \left(\frac{r}{a} \right)^4 \right)$$

C_e – adjusted iteratively to keep prescribed confinement



In addition, neoclassical ion heat conductivity

$$\chi_i = \chi_i^{an} + \chi_i^{neo}$$

Anomalous pinch $V_{pinch}/D_i \sim r/a^2$

Impurity ions

- Different types of impurities are treated simultaneously and self-consistently

$$\frac{\partial n_j^k}{\partial t} + \frac{1}{rg_1} \frac{\partial}{\partial r} (rg_2 \Gamma_j^k) = n_e \left[n_{j-1}^k \alpha_{ion,k}^{j-1} - n_j^k (\alpha_{ion,k}^j + \beta_{rec,k}^j) + n_{j+1}^k \beta_{rec,k}^{j+1} \right] \quad j = 1, \dots, Z_k$$

$$\Gamma_j^k = \Gamma_j^{nc,k} + \Gamma_j^{an,k}$$

High Z impurity accumulation

Pfirsch-Schlüter contribution

$$\Gamma_j^{nc} = -D_j^{PS,k} \partial n_j^k / \partial r + n_j^k W_j^{PS,k} = (1 + q^2) \rho_k^2 v_j^k \left[-\frac{\partial n_j^k}{\partial r} + Z_j \left(\frac{1}{n_i} \frac{\partial n_i}{\partial r} - \frac{1}{2T_i} \frac{\partial T_i}{\partial r} \right) \right]$$

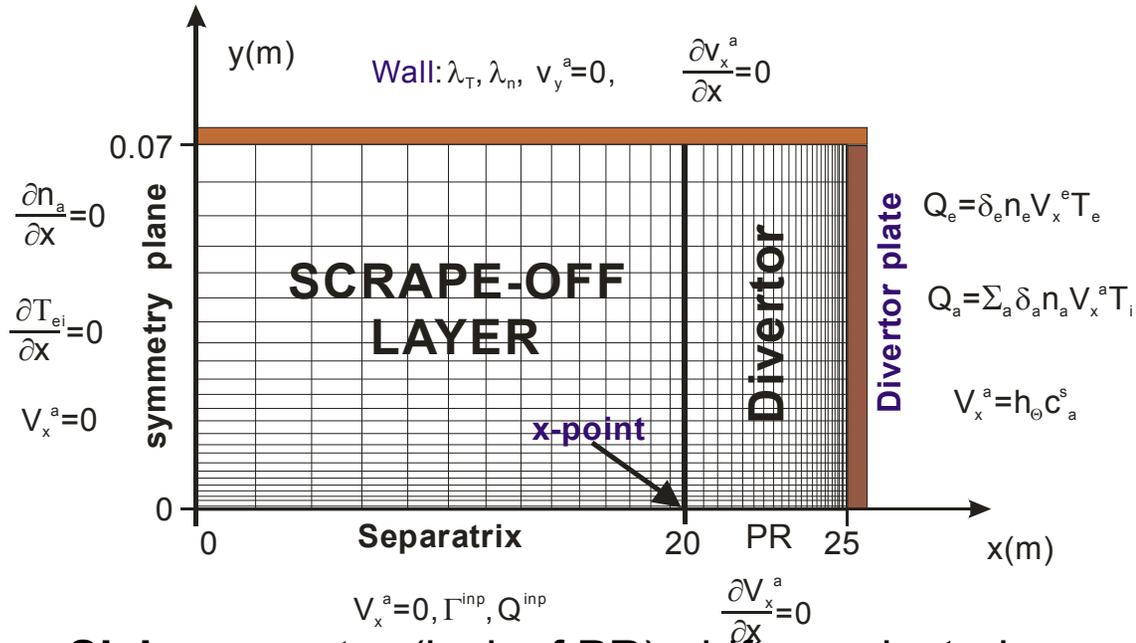
Anomalous contribution

$$\Gamma_j^{an,k} = -D_j^{an,k} \partial n_j^k / \partial r + n_j^k V_j^{pinch,k} \quad V_j^{pinch,k} \propto -D_j^{an,k} r / a^2$$

Usually anomalous transport same as for background plasma (ambipolarity)

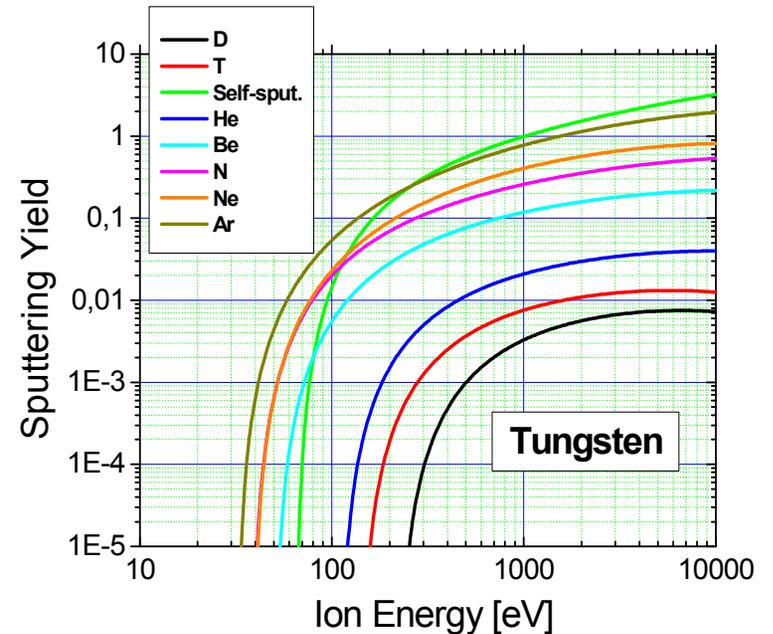
$$D_j^{an,k} = D_i^{an}$$

SOL Model



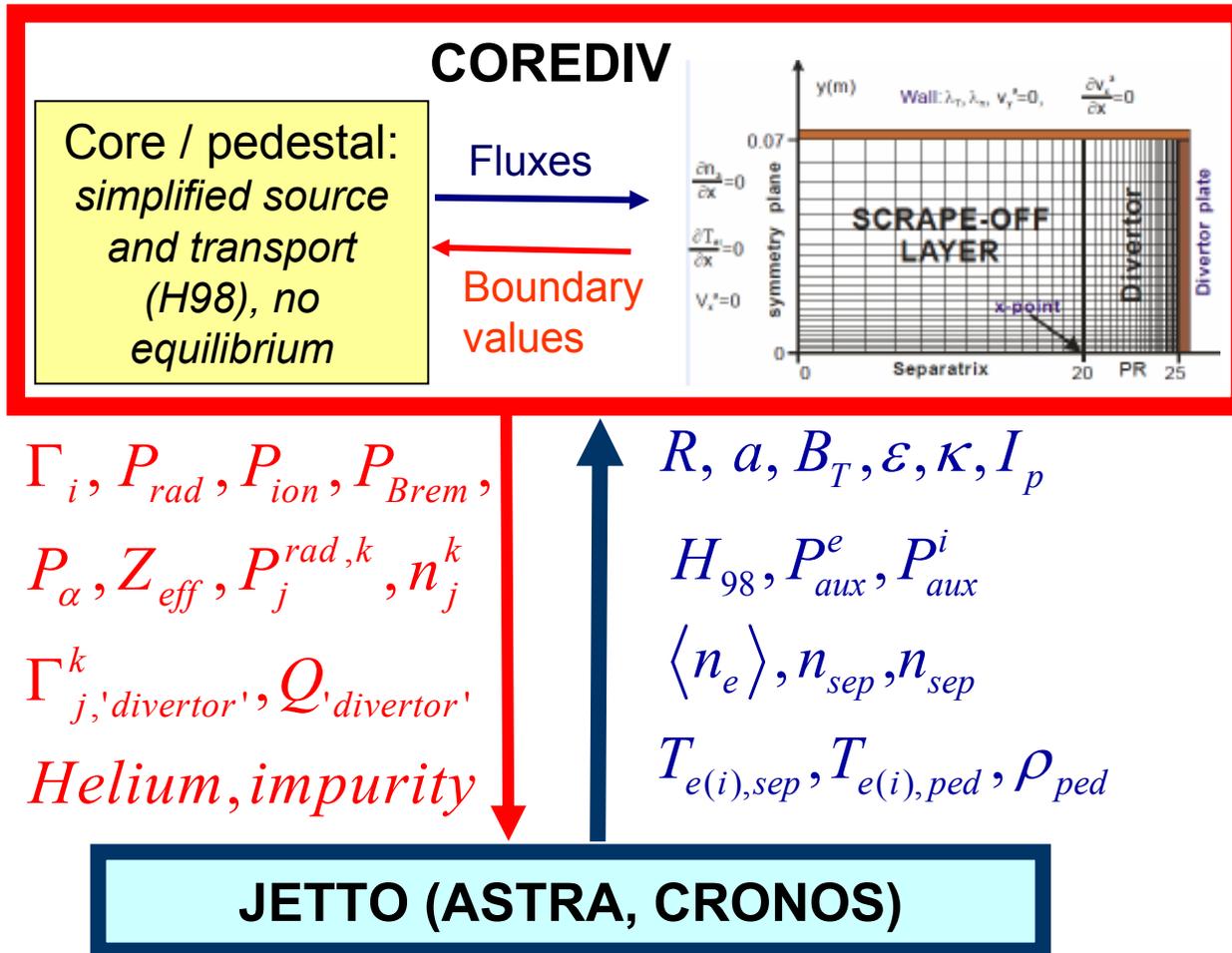
- **Slab geometry** (lack of PR), drifts neglected
- **Atomic processes**: ionization, recombination, excitation, charge exchange
- **Analytical model for neutrals** accounts for plasma recycling and impurity sputtering (also by seeded impurities). Recycling is an external parameter
- **Boundary conditions**: sheath, decay lengths; input fluxes from core part of the model
- **Intrinsic and seeded impurities** – gas puff at different positions

- 2D multifluid transport based on **Braginskii equations**
- Particle balance, parallel momentum, two energy equations
- Transport: parallel - classical, radial – anomalous



He, Li, Be, B, C, N, O, Ne,
Si, Ar, Ti, Ni, Mo, W

Coupled JETTO-COREDIV simulations for H-mode ITER plasmas



➤ ITER H-mode scenario: 15 MA, 5.34 T, 33 MW (NBI) + 20 MW (ECRH)

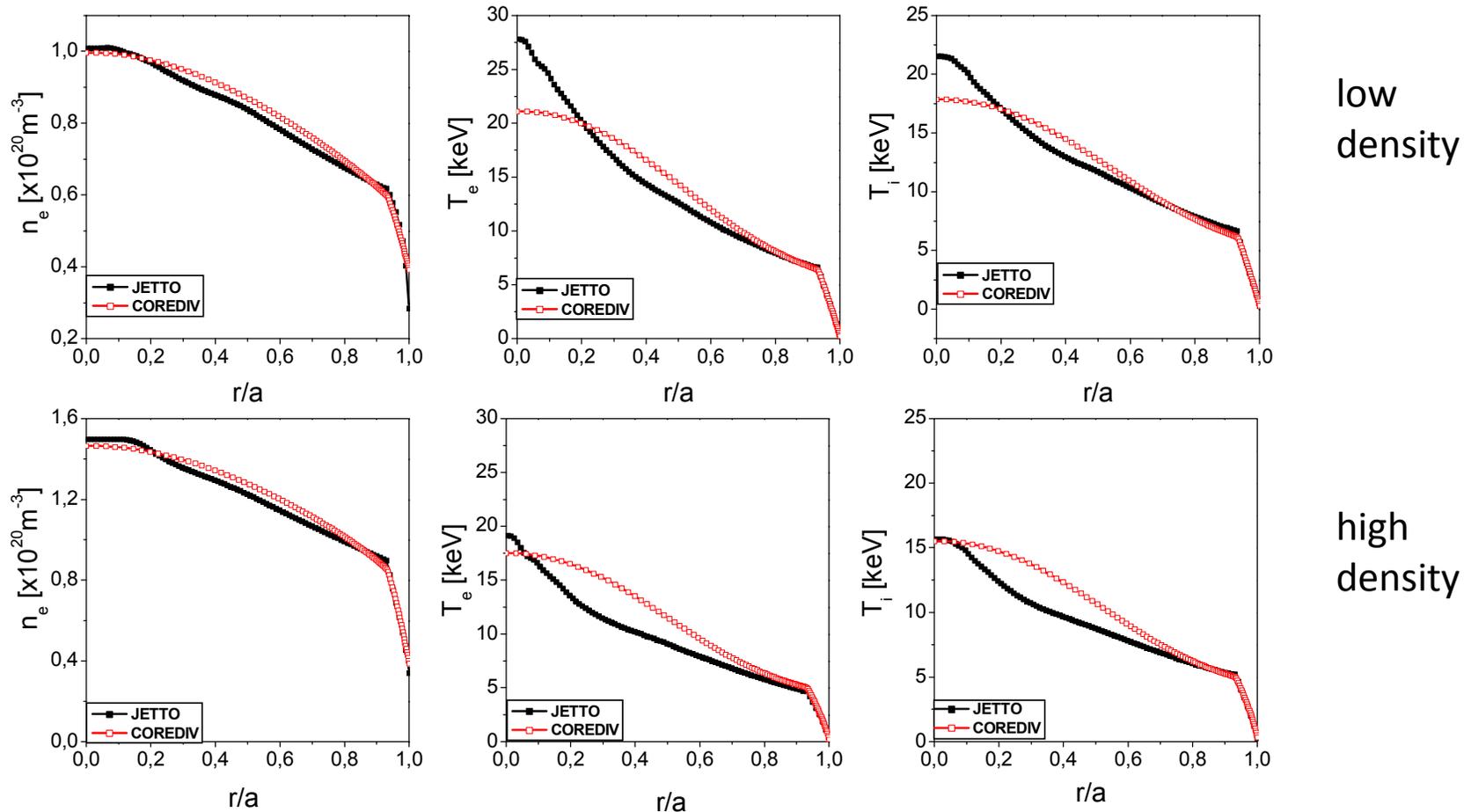
➤ JETTO / GLF23 (H98=0.85) with EPED and SOLPS boundary

➤ Low density:
 $n_{e_ped} = 6e19 \text{ m}^{-3}$
 $T_{e,ped} = 6.7 \text{ keV}$
 $T_{i,ped} = 7.5 \text{ keV}$

➤ High density:
 $n_{e_ped} = 9.5e19 \text{ m}^{-3}$
 $T_{e,ped} = 4.35 \text{ keV}$
 $T_{i,ped} = 4.83 \text{ keV}$

In ITER simulations we use the same version of COREDIV, which is used for JET!

Profile of n_e , T_e and T_i



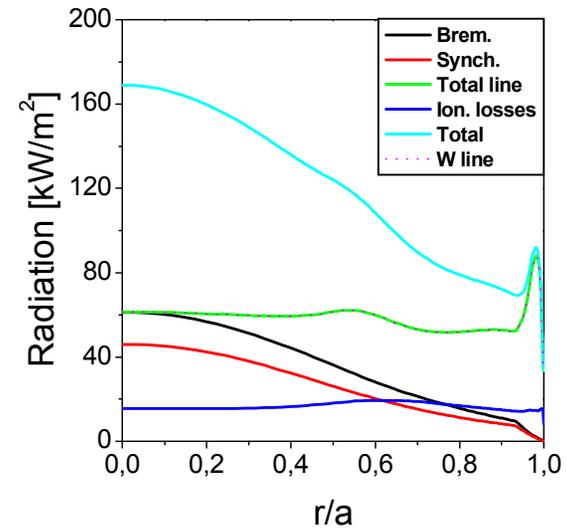
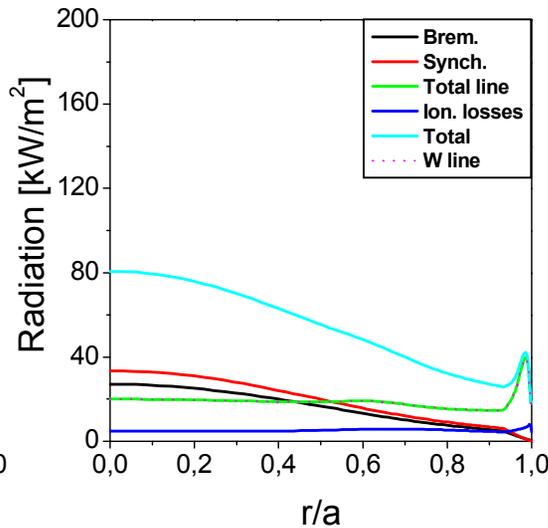
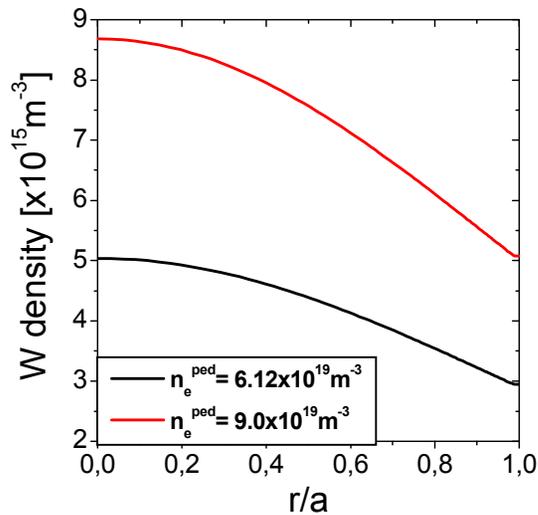
Comparison with JET simulation:

- high peaking of the density profile (4 times larger V_{pinch} than for JET plasmas)
- smaller edge barrier (for JET, $b = 0.15$, for ITER $b = 0.11$)

Profile of the tungsten density, radiation for two cases ($V_{\text{pinch}}^{\text{IMP}} = V_{\text{pinch}}^{\text{MAIN PLASMA}}$, without Ne seeding)

Case 1: low density

Case 2: high density



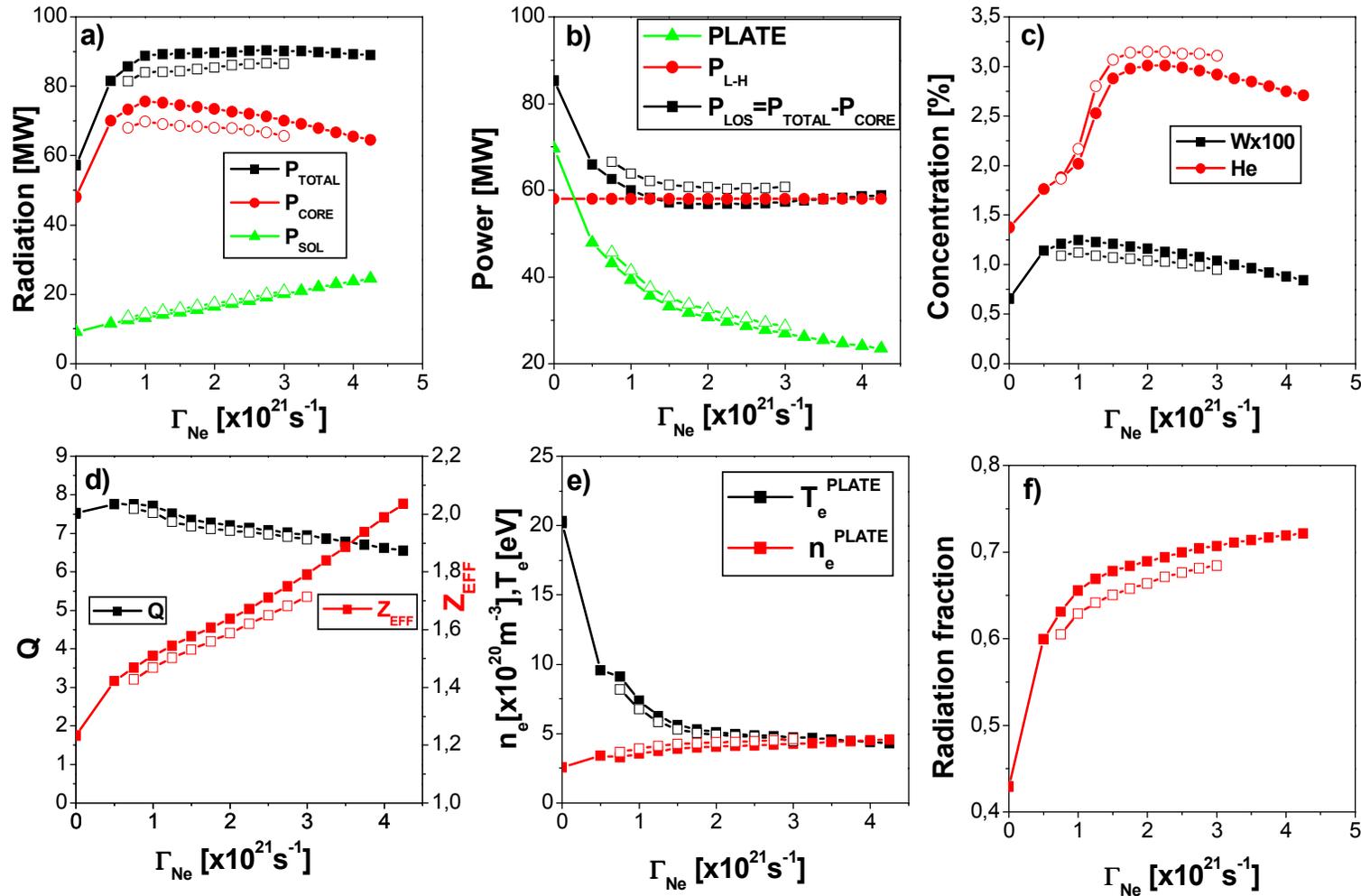
Parameters	Case 1	Case 2
P_{α} [MW]	79.4	140.1
P_{SYNCH} [MW]	11.67	15.1
P_{LIN} [MW]	19.92	42.52
P_{BREM} [MW]	9.71	19.98
P_{SOL} [MW]	83.65	80.4
P_{CORE} [MW]	41.92	89.97
Radiation Fraction	0.378	0.505
Total heating [MW]	133	193
Power to Plate [MW]	76.05	85.32
T_e^{PLATE} [eV]	25.45	33.55
n_e^{PLATE} [$\times 10^{20} \text{m}^{-3}$]	2.1	1.43
Recycling coefficient	0.9958	0.989

Power to plate is not acceptable for operation window for ITER divertor (10 MW/m^2).

ITER needs gas puffing!

Sputtering	Case 1	Case 2
D	3.17	4.42
W(self-sputtering)	7.806	14.505
He	0.781	0.502

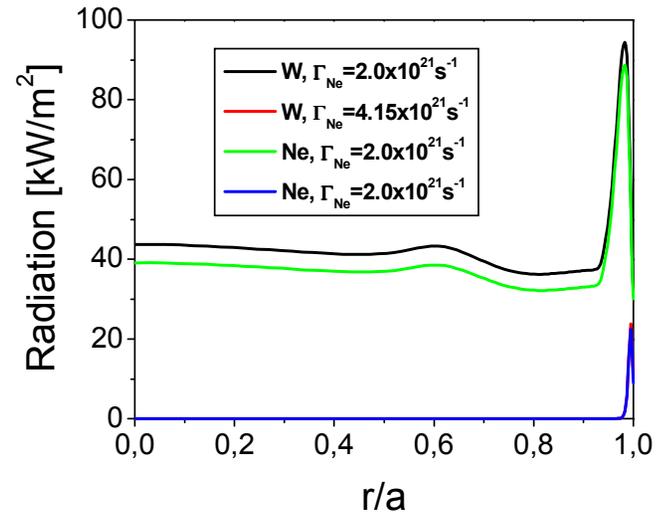
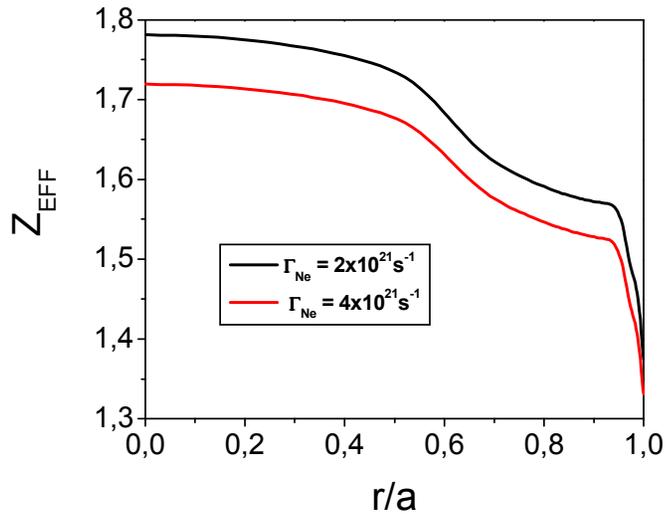
Ne gas puff for two different density on the separatrix



$n_{sep} = 3.25 \times 10^{19} \text{ m}^{-3}$ (full symbols)
 $n_{sep} = 3.5 \times 10^{19} \text{ m}^{-3}$ (open symbols)

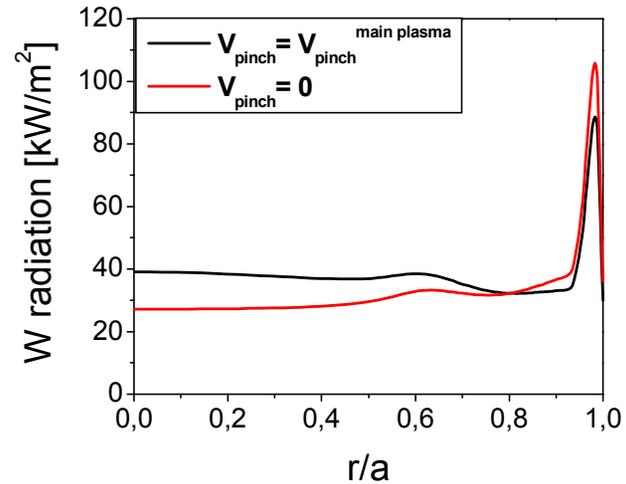
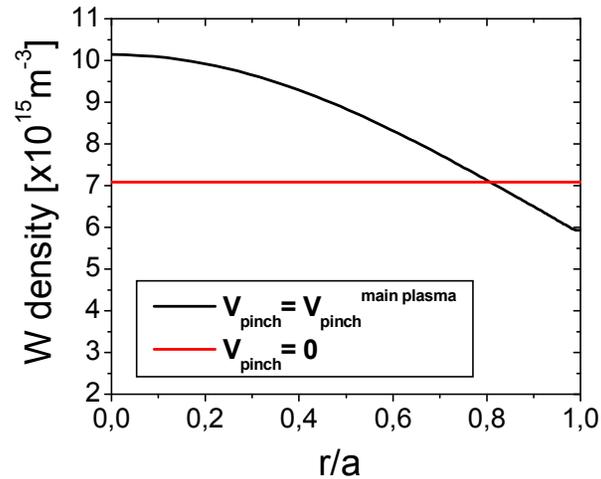
The effect of separatrix density on the power through the separatrix is weak!

Influence on the Ne gas puff to the profile on the Z_{EFF} and radiation



Tungsten density n_W (left) and Ne radiation (right) with $\Gamma_{\text{Ne}} = 2 \times 10^{21}$ and 4.15×10^{21} part/s

Effect on the $V_{\text{pinch}}^{\text{IMP}}$



*Tungsten density and radiation obtained in simulations
with $V_W = V_D$ (black) and $V_W = 0$ (red) with $\Gamma_{\text{Ne}} = 2 \times 10^{21} \text{ part/s}$*

- for zero W inward pinch $V_W = 0$ the W concentration reduces only by 8.5%
- core radiation reduces only by 2.4 MW

CONCLUSION

Taking into account two essential constraints:

- ❖ H-mode operation (i.e. $P_{\text{loss}} > P_{\text{LH}}$)
- ❖ low divertor heat loads (power to divertor does not exceed 40 MW)
 - Ne seeding is essential for reducing the power to plate below 40 MW
 - W sputtering by Ne is important, it replaces the W self-sputtering leading to a larger W accumulation and core radiation, than in the case without Ne seeding
 - H-mode operational point with power to plate slightly below 40 MW is barely above the L-H power threshold in regimes with Ne seeding due to large W radiation. The H-mode operation is sensitive to the separatrix density and W inward pinch.

It should be mentioned that the strong coupling between the W accumulation (or transport) and radiation, reducing the power to plate, divertor temperature and W production, which in its turn affects the W accumulation leads to a “stiff” operational point and make it difficult to estimate in advance the trend with engineering parameters. More extensive parameter scans are needed to determine if the operational space for robust H-mode ITER performance with high fusion yield and acceptable level of divertor heat loads exist.